

The prime symbol denotes exclusion from the sum of the quantities having the same subscripts  $\nu$  and  $n$ . The downwash factor  $\omega$  is defined by the relation

$$\omega = 2n \quad (10)$$

Reference 1 takes an average value of  $\omega$  equal to  $2n_s$ , which is the sheared value for a swept wing [ $\lambda(y) = 0$ ]. For a  $45^\circ$  swept wing with aspect ratios of 3, 5, and 7, the value for  $\omega$  is 1.025, 1.010, and 1.005, respectively. Strictly speaking,  $\omega$  is also dependent on the spanwise location and would vary in the same manner as  $n$  varies, since  $\omega = 2n$ .

The load distribution obtained by using Eq. (9), for  $m = 31$ , is shown in Figs. 2a, 2b, 2c, and 2d for the wing planforms of Refs. 4, 5, 6, and 7, respectively. In accordance with Refs. 1 and 10, the value for  $\omega$  was held constant; a value of unity was assigned following Ref. 10. Figure 2 shows the Kuchemann-Multhopp lifting line method using aerodynamic centers obtained by the methods of Refs. 1, 2, and 8 on the upper set of curves in each figure. The difference in load distribution between that calculated using the Kuchemann tangent approximation and the Kuchemann hyperbolic approximation is small, so that the loading curves shown make use of the tangent approximation. In addition, no consideration was given to wing thickness effects, except in its effect on the aerodynamic center position in using the method of Ref. 2, in the theoretical calculations so the value of  $a_0$  used was  $2\pi$ . In general, of the loadings obtained by the Kuchemann-Multhopp lifting line method, that obtained from using the aerodynamic centers from the Kuchemann tangent approximation appears to give the best agreement with experiment except for the case of wing D. The loading for wing D appears to be best represented by use of either the aerodynamic centers from the Multhopp lifting surface theory<sup>8</sup> or the Transonic Data Memorandum method<sup>2</sup> without the thickness correction. The dotted curves shown in Fig. 2c show the slight reduction in the calculated loading obtained by using the extended aerodynamic center distribution shown by the dashed lines in Fig. 1c.

Figure 2 also shows the span loading characteristics of wings A, B, C, and D as obtained from the Weissinger<sup>11</sup> (using 59 spanwise points) and Multhopp lifting surface<sup>8</sup> (using 4 chordwise and 31 spanwise points) theories compared with experiment. This figure shows that the loadings predicted by the Weissinger theory and Multhopp lifting surface theory give comparable results and agree quite well with experiment except in the case of wing D where both theories slightly underpredict the level of the loading on the wing.

### References

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## Change in Pitching-Moment Coefficient Due to Ground Effect

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### Nomenclature

- $a_t$  = slope of tail lift coefficient  
 $C_L$  = lift coefficient  
 $C_m$  = pitching-moment coefficient  
 $\bar{c}$  = mean aerodynamic chord  
 $q$  = dynamic pressure  
 $S$  = wing area  
 $\alpha$  = angle of attack  
 $\Delta$  = denoting the change due to ground effect in case of constant angle of attack, e.g.,  $\Delta C_L = C_L - C_{L_0}$   
 $\delta_e$  = elevator angle  
 $\epsilon$  = downwash angle at the tail  
 $\tau$  = relative control effectiveness,  $\tau = \partial a_t / \partial \delta_e$

### Subscripts

- $t$  = tail  
 $w$  = wing-body combination  
 $0$  = free air, e.g.,  $C_{L_0}$   
 $|\delta$  = constant elevator angle

REFERRING to a new flight-test method for measurement of ground effect of fixed wing aircraft,<sup>1</sup> this Note is concerned with the change in pitching-moment coefficient due to ground effect and the factors causing this change in case of a constant-angle-of-attack approach. The investigation described herein has been derived from a method which was used for evaluating the change in downwash angle at the tail when measuring ground effect of the Transall C-160 airplane.<sup>2</sup>

Dividing up the forces and moments as shown in Fig. 1, the change in lift coefficient due to ground effect can be expressed as

$$\Delta C_L = \Delta C_{L_w} + k_t(\Delta C_{L_t}|\delta + \tau a_t \Delta \delta_e) \quad (1)$$

where terms of small magnitude are neglected and

$$k_t = S_t q_{t_0} / S q_0 = S_t (q_{t_0} + \Delta q_t) / [S (q_0 + \Delta q)] \quad (2)$$

is assumed to be constant. The difference between the moment equations in free air and in ground proximity yields

$$x_a \Delta C_{L_w} - k_t(l'_t - x_a)(\Delta C_{L_t}|\delta + \tau a_t \Delta \delta_e) = 0 \quad (3)$$

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